



**University of
Zurich**^{UZH}

**Zurich Open Repository and
Archive**

University of Zurich
University Library
Strickhofstrasse 39
CH-8057 Zurich
www.zora.uzh.ch

Year: 2010

Top quark phenomenology

Rikkert, F

Abstract: In this talk three 2-sigma deviations from the Standard Model predictions in the top quark sector are briefly discussed. These are the excess of events in the tail of the H_T distribution in $t\bar{t}$ events, the top-quark charge asymmetry and the discrimination of s- and t-channel events in single top. The latter has only been observed by CDF, while the other two are found by both CDF and D0.

Posted at the Zurich Open Repository and Archive, University of Zurich
ZORA URL: <https://doi.org/10.5167/uzh-41856>
Conference or Workshop Item
Accepted Version

Originally published at:

Rikkert, F (2010). Top quark phenomenology. In: Hadron Collider Physics Symposium 2010: HCP 2010, Toronto, Ontario, Canada, 23 August 2010 - 27 August 2010.

Top Quark Phenomenology

Rikkert Frederix

Institute for Theoretical Physics, University of Zurich, Zurich, Switzerland

In this talk three 2-sigma deviations from the Standard Model predictions in the top quark sector are briefly discussed. These are the excess of events in the tail of the H_T distribution in $t\bar{t}$ events, the top-quark charge asymmetry and the discrimination of s- and t-channel events in single top. The latter has only been observed by CDF, while the other two are found by both CDF and DØ.

I. TOP QUARKS AT THE TEVATRON

Everything we know about the top quark is coming from the experiments at the Tevatron collider at Fermilab. Since its discovery in 1995 [1, 2], many of its properties have been firmly established over the years. The most precise measurements are its mass [3] and cross section [4, 5], but results for other observables and properties of the top quark are improving as well. This includes: the top pair invariant mass distribution [6, 7], the forward-backward charge asymmetry [8, 9], the top quark width and its lifetime [10, 11], the branching fraction to Wb compared to Wq [12, 13], the W -boson helicity fractions [14, 15], spin correlation in the top quark pair production [14, 16], the top pair production mechanism (gg versus $q\bar{q}$ initial state) [17], the top quark's electric charge [18, 19], the mass difference between top and anti-top quarks [20, 21], the total single top cross section [22] including bounds on the CKM matrix element $|V_{tb}|$ [22] and discrimination between the single-top s- and t-channel production mechanisms [23, 24]. Furthermore there have been many searches for new physics in the top quark sector, such as: resonant contributions to the $t\bar{t}$ production [6, 7], search for fourth generation t' and b' quarks [25, 26], scalar top admixture in $t\bar{t}$ events [27], top decays to charged Higgs bosons [28, 29], W' and H^+ resonant contributions to s-channel single top [30–32], anomalous Wtb couplings [33], production via a FCNC ($u(c) + g \rightarrow t$) [34, 35] and decay in $t \rightarrow Zq$ [36].

All measurements are in agreement with the Standard Model expectations within uncertainties. There are, however, three measurements that are slightly off: in the search for t' quarks, the tail of the H_T distribution has a small excess of events, the top quark charge asymmetry is more pronounced than expected and the discrimination of the s- and t-channel cross section in single top is off by more than 2 sigma (CDF only). In the rest of this talk I'll give my personal view on these three experimental results.

II. THE THREE 2-SIGMA DEVIATIONS

A. The Tail Of The H_T Distribution

The H_T observable is defined as the scalar sum of all transverse energies of the jets, leptons and missing E_T and is a measure for the overall scale of the process. Recently both CDF and DØ have updated their analysis for the search of heavy fourth generation quarks, in particular in t' , i.e. a fourth generation top quark, that decays predominantly to W^+b . After selection cuts, the H_T distribution shows an excess of events in the tail of the distribution. In fig. 1(a), this distribution is shown for DØ. As can be clearly seen from the last two bins (where the final bin is an overflow bin and contains also all events with $H_T > 700$ GeV) there are slightly more events than could have been expected from $t\bar{t}$, W/Z + jets and multi-jet backgrounds. The inclusion of a heavy t' quark to the physics model, increases the number of events in the tail of the distribution and results in a better description of the data. This can also be seen from the plot in fig 1(b), where for a t' mass larger than about 350 GeV the observed limit is at the border of the 95% confidence level for the predicted limit. This means that limit on the cross section for t' pairs is much milder than expected, which reflects the slight excess of events found in the tail of the H_T distribution.

From a theoretical point of view, the tail of the H_T distribution is difficult to model, because it is an observable quite sensitive to higher order corrections, see e.g. Ref. [37]. Furthermore, the approximation of describing extra

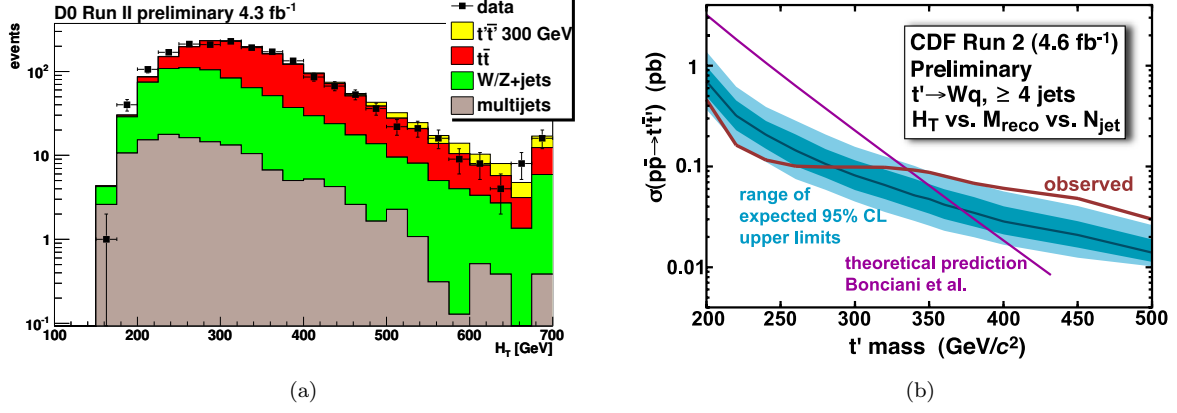


FIG. 1: The H_T distribution as measured by the DØ collaboration (a). Observed (red curve) and expected (blue band) limits on the t' pair production cross section as a function of the t' mass as extracted from the data by CDF are shown (b). Figures are taken from Ref. [26] and [25], respectively.

radiation by a logarithm, as is done in all parton showers, does not capture hard radiation correctly, see e.g. Ref. [38]. Using a CKKW [39] or MLM [40] technique to merge higher multiplicity matrix elements with the parton shower could help in describing the process. This has been done for the $W + \text{jets}$ background, but not for the $t\bar{t}$ final state and therefore this might have an impact on the predictions and increase the number $t\bar{t}$ events in the tail and explain the slight excess.

Given that the effect is small (2 sigma) and non-trivial to model by Monte Carlo, it is not (yet) significant. More data, and in particular the LHC with its larger center-of-mass energy, and more refined analyses using matrix-elements with parton-shower matching will shed light on these effects.

B. Top Quark Charge Asymmetry

In pure QCD production of top quark pairs at hadron colliders, the lowest order prediction shows no preferred direction for the top and anti-top quarks. When the next-to-leading corrections are taken into account, the interference between gluon radiation from the initial state quark line and the top quark line in the $q\bar{q} \rightarrow t\bar{t}g$ channel has a negative contribution to the charge asymmetry, while the contribution from the interference between the virtual box diagrams and the Born is positive. Furthermore there is a negligible contribution to the charge asymmetry from the flavor excitation processes $qg \rightarrow t\bar{t}q$. Top quark production by gluon fusion is symmetric in its initial state and therefore does not contribute to the asymmetry.

Quantitatively, the contribution from the virtual corrections is larger, hence the top quarks prefer to go in the direction of the incoming quark, and the anti-top quarks in the directions of the incoming anti-quark. Including the (LO) electroweak effects for the Tevatron the charge asymmetry is given by [41, 42]

$$\begin{aligned} A_{FB}(\text{lab}) &= 0.051 \pm 0.006, \\ A_{FB}(t\bar{t}) &= 0.078 \pm 0.009, \end{aligned} \tag{1}$$

where the uncertainties are coming from renormalization and factorization scale dependence and

$$\begin{aligned} A_{FB}(\text{lab}) &= \frac{\int_{y>0} N_t(y) - \int_{y>0} N_{\bar{t}}(y)}{\int_{y>0} N_t(y) + \int_{y>0} N_{\bar{t}}(y)} \\ A_{FB}(t\bar{t}) &= \frac{\int N(\Delta y > 0) - \int N(\Delta y < 0)}{\int N(\Delta y > 0) + \int N(\Delta y < 0)}, \end{aligned} \tag{2}$$

are the definition of the charge asymmetry in the lab and the $t\bar{t}$ rest frames, respectively. Here, $N_{t(\bar{t})}(y)$ is the number of top (anti-top) quarks as a function of the rapidity y and $\Delta y = y_t - y_{\bar{t}}$ the difference in rapidities between the top and anti-top quarks.

Recently, the prediction for the asymmetry in the top pair rest frame has been improved by including threshold logarithms at all orders and it was found that the results are quite stable under inclusion of these all-order effects [43, 44],

$$A_{FB}(t\bar{t}) = 0.073^{+0.009}_{-0.007}. \quad (3)$$

Both the CDF and DØ collaborations measure a non-zero top quark charge asymmetry. CDF finds the following values [8]

$$\begin{aligned} A_{FB}(\text{lab}) &= 0.073 \pm 0.028 \quad \text{CDF} \\ A_{FB}(t\bar{t}) &= 0.057 \pm 0.028 \quad \text{CDF}, \end{aligned} \quad (4)$$

while DØ has only performed the analysis for the asymmetry in the top pair rest frame [9]

$$A_{FB}(t\bar{t}) = 0.08 \pm 0.04 \text{ (stat.)} \pm 0.01 \text{ (syst.)} \quad \text{DØ}. \quad (5)$$

However, these results are uncorrected for hadronization, underlying event, background effects, etc. which makes a direct comparison with the predictions, eq. (1) and (3), unfair. The CDF collaboration provides also results corrected for these effects [8], which enhances the the asymmetry,

$$\begin{aligned} A_{FB}(\text{lab}) &= 0.150 \pm 0.050 \text{ (stat.)} \pm 0.024 \text{ (syst.)} \quad \text{CDF (corrected)} \\ A_{FB}(t\bar{t}) &= 0.158 \pm 0.072 \text{ (stat.)} \pm 0.017 \text{ (syst.)} \quad \text{CDF (corrected)}. \end{aligned} \quad (6)$$

These measurements are almost two standard deviations larger than the theory predictions given in eq. (1) and (3). Notice that the uncorrected results from CDF and DO are in agreement with each other, and if we assume that the unfolding to parton level is the same for DØ as it is for CDF, also the asymmetry as measured by DØ is also about 2 sigma from its predicted value in perturbative QCD.

Even though the effect is around two standard deviations, and could therefore very well be a statistical effect, from the theory point of view many models beyond the SM were studied to see if these effects could be explained by new physics, see e.g. Ref. [45] and references therein. However, before assigning the large measured asymmetry to BSM effects, another option should be considered as well. Even though the NLO (i.e. first non-zero) predictions are stable under all-order resummation of threshold logarithms, this does not take into account that there might be a sizable effect from higher order virtual corrections, i.e. the two loop corrections to $q\bar{q} \rightarrow t\bar{t}$. The need to go from NLO to the full NNLO is acknowledged.

The ingredients to go one full order higher in perturbation theory are almost all known in analytic form [46–49] and the remaining contributions can be computed numerically [50]. The one important missing ingredient is the prescription to combine the contributions in an infra-red finite way. Several methods have been proposed [51–54], but non of them have been proven to work for top pair production in hadron collisions so far. Obtaining these results needs further theoretical developments.

From the experimental point of view the measurements are still statistically dominated, so with more data the measurements will become more precise. Due to the fact that the LHC as a proton-proton collider and that the initial state is dominated by gluons, measuring the asymmetry at this collider will be challenging at the least.

C. Discrimination Between s- And t-channel Events By CDF

Single top quark production has been firmly established with 5 standard deviations last year [55, 56]. At the Tevatron only two of the three single top production channels contribute significantly, i.e. the s channel (with a cross

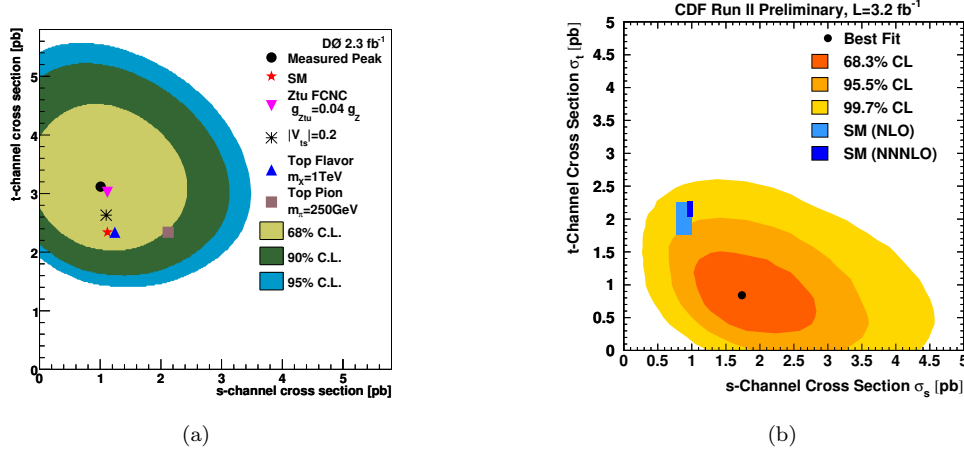


FIG. 2: Posterior probability for s- and t-channel single top production by the DØ (a) and CDF (b) collaborations. Figures are taken from Ref. [24] and [23], respectively.

section around 1 pb) and the t channel (about 2 pb). The total single top rate measured at the Tevatron agrees very well with the sum of the s- and t-channel cross section predictions.

For DØ also the discrimination between the s- and t-channel cross section agrees well within 1 sigma with the theoretical predictions, see fig. 2(a). For CDF the situation is worse: the theoretical predictions are just outside the 95.5% C.L. contour, as can be seen in fig. 2(b). This discrimination is made using multi-variate techniques that use as much information as possible from the event topologies [23, 24].

It is known that the t-channel process, see fig. 3(a), is difficult to model by Monte Carlo event generators due to the initial state b quark. After showering of the events, the initial state b quark is modeled (in most cases) to come from a gluon splitting into a $b\bar{b}$ pair, for which the \bar{b} quark becomes a final state parton, which is called the spectator- b in the following. When only using the $2 \rightarrow 2$ process as in fig. 3(a), the transverse momentum of the spectator- b is greatly underestimated by the parton shower. This has been acknowledged and addressed by applying a merging procedure for leading order $2 \rightarrow 2$ and $2 \rightarrow 3$ events [57]. There is much freedom in the precise choice of the merging which can lead to a large spread in predictions for the spectator- b .

Recently, the NLO corrections to t-channel single top production in the 4-flavor scheme, i.e. starting from the $2 \rightarrow 3$ process, fig. 3(b), were performed [58, 59]. This calculation gives the distributions for the spectator- b quark for the first time at NLO, and should therefore be considered the predictions to validate the merging procedure of the $2 \rightarrow 2$ and $2 \rightarrow 3$ MC samples against.

Comparing the new predictions for the fraction of spectator- b quarks with large transverse momentum ($p_T > 20$ GeV) and small pseudo rapidities ($|\eta| < 2$) it has been found that results from the matching procedure as used by the DØ collaboration follows the NLO predictions quite well [60]. Unfortunately, for the Monte Carlo samples used by the CDF collaboration, which are based on the ZTOP program [61], this ratio is about half the size as the NLO prediction [62]. This means that CDF assumed that the t-channel events have only about half as many spectator- b

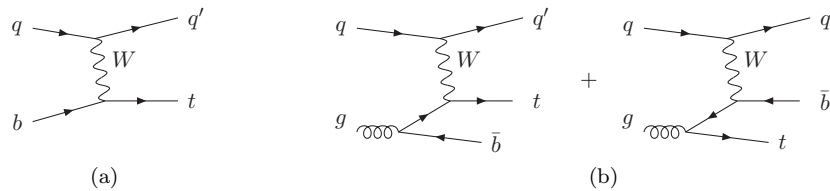


FIG. 3: Lowest order Feynman diagrams for the t-channel single-top process in the five (a) and four (b) flavor schemes.

quarks as predicted by NLO. Given that s-channel events have an extra b quark compared to the $2 \rightarrow 2$ t-channel events, at first sight this discrepancy could explain the measured excess of s-channel events compared to t-channel events: many t-channel events with a spectator- b quark in the data were assigned to belong to the s-channel cross section. However, a careful improvement in the analysis by the CDF collaboration taking into account the new NLO computation, shows that this effect on the discrimination of the s- and t-channel cross sections is negligibly small [62]. Therefore, even with the improved theoretical predictions the 2 standard deviations discrepancy remains.

If the 2-sigma effect persists or increases when more data is collected and analysed by CDF, experimental effort to re-address the estimation of the uncertainties might be acknowledged. Also more theoretical work is needed to further improve predictions, e.g. by matching the NLO t-channel single top calculation in the 4-flavor scheme to a parton shower program following e.g. the MC@NLO [63] or POWHEG [64] methods.

III. CONCLUSIONS

There are three measurements in the top quark sector that are about 2 standard deviations away from its predictions within the Standard Model:

- There is a small excess of events in the tail of the H_T distribution. Possible explanations might be that for the H_T observable the tail of the distribution is known to be non-trivial to model using Monte Carlo event generators. In particular, the uncertainty from higher order corrections can be large. However, if the effect is truly there, a possible explanation is a fourth generation of quarks.
- The top quark charge asymmetry is found to be more pronounced than predicted by NLO QCD, which is the first order to give a non-zero result. The question arises if higher order effects might be important here, and it has been shown by using resummation techniques that this is probably not the case for the (soft) real emission. However, given that at NLO the virtual corrections give a larger contribution than the real emission, also the virtual corrections should be calculated at higher order to provide a prediction at high enough precision.
- The discrimination of s- and t-channel events in single top production by CDF. Even though in the Monte Carlo predictions used in the CDF analyses there were half as many spectator- b quarks as compared to the recent NLO calculation, it does not explain the difference with $D\bar{O}$ nor with the theoretical predictions. There is no understanding yet of this peculiarity.

On the other hand, these three 2-sigma deviations could very well be statistical fluctuations. It is therefore important to see what this will give with more data, and what the LHC will tell is in the near future.

-
- [1] **CDF** Collaboration, F. Abe *et al.*, Phys. Rev. Lett. **74**, 2626 (1995), [hep-ex/9503002].
 - [2] **DØ** Collaboration, S. Abachi *et al.*, Phys. Rev. Lett. **74**, 2632 (1995), [hep-ex/9503003].
 - [3] **CDF & DØ** Collaboration, arXiv:1007.3178 [hep-ex].
 - [4] **CDF** Collaboration, Conf. Note 9913 (2009).
 - [5] **DØ** Collaboration, Conf. Note 6038 (2010).
 - [6] **CDF** Collaboration, T. Aaltonen *et al.*, arXiv:0903.2850 [hep-ex].
 - [7] **DØ** Collaboration, Conf. Note 5882 (2009).
 - [8] **CDF** Collaboration, Conf. Note 10185 (2010).
 - [9] **DØ** Collaboration, Conf. Note 6062 (2010).
 - [10] **CDF** Collaboration, Conf. Note 10035 (2010).
 - [11] **DØ** Collaboration, Conf. Note 6034 (2010).
 - [12] **CDF** Collaboration, D. E. Acosta *et al.*, Phys. Rev. Lett. **95**, 102002 (2005), [hep-ex/0505091].
 - [13] **DØ** Collaboration, V. M. Abazov *et al.*, Phys. Rev. Lett. **100**, 192003 (2008), [arXiv:0801.1326 [hep-ex]].

- [14] **CDF** Collaboration, Conf. Note 10211 (2010).
- [15] **DØ** Collaboration, Conf. Note 5722 (2008).
- [16] **DØ** Collaboration, Conf. Note 5950 (2009).
- [17] **CDF** Collaboration, T. Aaltonen *et al.*, Phys.Rev. **D78**, 111101 (2008), [arXiv:0712.3273 [hep-ex]].
- [18] **CDF** Collaboration, Conf. Note 9939 (2010).
- [19] **DØ** Collaboration, V. Abazov *et al.*, Phys.Rev.Lett. **98**, 041801 (2007), [hep-ex/0608044].
- [20] **CDF** Collaboration, Conf. Note 10173 (2010).
- [21] **DØ** Collaboration, V. Abazov *et al.*, Phys.Rev.Lett. **103**, 132001 (2009), [arXiv:0906.1172 [hep-ex]].
- [22] **CDF & DØ** Collaboration, T. E. W. Group, arXiv:0908.2171 [hep-ex].
- [23] **CDF** Collaboration, T. Aaltonen *et al.*, arXiv:1004.1181 [hep-ex].
- [24] **DØ** Collaboration, V. M. Abazov *et al.*, Phys.Lett. **B682**, 363 (2010), [arXiv:0907.4259 [hep-ex]].
- [25] **CDF** Collaboration, Conf. Note 10110 (2010).
- [26] **DØ** Collaboration, Conf. Note 5892 (2010).
- [27] **DØ** Collaboration, V. Abazov *et al.*, Phys.Lett. **B674**, 4 (2009), [arXiv:0901.1063 [hep-ex]].
- [28] **CDF** Collaboration, Conf. Note 10104 (2010).
- [29] **DØ** Collaboration, V. Abazov *et al.*, Phys.Lett. **B682**, 278 (2009), [arXiv:0908.1811 [hep-ex]].
- [30] **CDF** Collaboration, T. Aaltonen *et al.*, Phys.Rev.Lett. **103**, 041801 (2009), [arXiv:0902.3276 [hep-ex]].
- [31] **DØ** Collaboration, V. Abazov *et al.*, Phys.Rev.Lett. **100**, 211803 (2008), [arXiv:0803.3256 [hep-ex]].
- [32] **DØ** Collaboration, V. Abazov *et al.*, Phys.Rev.Lett. **102**, 191802 (2009), [arXiv:0807.0859 [hep-ex]].
- [33] **DØ** Collaboration, V. Abazov *et al.*, Phys.Rev.Lett. **102**, 092002 (2009), [arXiv:0901.0151 [hep-ex]].
- [34] **CDF** Collaboration, T. Aaltonen *et al.*, Phys.Rev.Lett. **102**, 151801 (2009), [arXiv:0812.3400 [hep-ex]].
- [35] **DØ** Collaboration, V. M. Abazov *et al.*, arXiv:1006.3575 [hep-ex].
- [36] **CDF** Collaboration, T. Aaltonen *et al.*, Phys.Rev.Lett. **101**, 192002 (2008), [arXiv:0805.2109 [hep-ex]].
- [37] M. Rubin, G. P. Salam and S. Sapeta, arXiv:1006.2144 [hep-ph].
- [38] J. Alwall, S. de Visscher and F. Maltoni, JHEP **0902**, 017 (2009), [arXiv:0810.5350 [hep-ph]].
- [39] S. Catani, F. Krauss, R. Kuhn and B. Webber, JHEP **0111**, 063 (2001), [hep-ph/0109231].
- [40] J. Alwall *et al.*, Eur.Phys.J. **C53**, 473 (2008), [arXiv:0706.2569 [hep-ph]].
- [41] J. H. Kuhn and G. Rodrigo, Phys.Rev.Lett. **81**, 49 (1998), [hep-ph/9802268].
- [42] J. H. Kuhn and G. Rodrigo, Phys.Rev. **D59**, 054017 (1999), [hep-ph/9807420].
- [43] L. G. Almeida, G. F. Sterman and W. Vogelsang, Phys.Rev. **D78**, 014008 (2008), [arXiv:0805.1885 [hep-ph]].
- [44] V. Ahrens, A. Ferroglia, M. Neubert, B. D. Pecjak and L. L. Yang, arXiv:1003.5827 [hep-ph].
- [45] Q.-H. Cao, D. McKeen, J. L. Rosner, G. Shaughnessy and C. E. M. Wagner, Phys. Rev. **D81**, 114004 (2010), [arXiv:1003.3461 [hep-ph]].
- [46] R. Bonciani, A. Ferroglia, T. Gehrmann, D. Maitre and C. Studerus, JHEP **0807**, 129 (2008), [arXiv:0806.2301 [hep-ph]].
- [47] R. Bonciani, A. Ferroglia, T. Gehrmann and C. Studerus, JHEP **0908**, 067 (2009), [arXiv:0906.3671 [hep-ph]].
- [48] J. G. Korner, Z. Merebashvili and M. Rogal, Phys. Rev. **D77**, 094011 (2008), [arXiv:0802.0106 [hep-ph]].
- [49] S. Dittmaier, P. Uwer and S. Weinzierl, Phys. Rev. Lett. **98**, 262002 (2007), [hep-ph/0703120].
- [50] M. Czakon, Phys.Lett. **B664**, 307 (2008), [arXiv:0803.1400 [hep-ph]].
- [51] A. Gehrmann-De Ridder, T. Gehrmann and E. W. N. Glover, JHEP **09**, 056 (2005), [hep-ph/0505111].
- [52] S. Catani and M. Grazzini, Phys. Rev. Lett. **98**, 222002 (2007), [hep-ph/0703012].
- [53] M. Czakon, Phys.Lett. **B693**, 259 (2010), [arXiv:1005.0274 [hep-ph]].
- [54] G. Somogyi, Z. Trocsanyi and V. Del Duca, JHEP **0506**, 024 (2005), [hep-ph/0502226].
- [55] **DØ** Collaboration, V. M. Abazov, arXiv:0903.0850 [hep-ex].
- [56] **CDF** Collaboration, T. Aaltonen *et al.*, arXiv:0903.0885 [hep-ex].
- [57] E. E. Boos, V. E. Bunichev, L. V. Dudko, V. I. Savrin and A. V. Sherstnev, Phys. Atom. Nucl. **69**, 1317 (2006).
- [58] J. M. Campbell, R. Frederix, F. Maltoni and F. Tramontano, JHEP **10**, 042 (2009), [arXiv:0907.3933 [hep-ph]].
- [59] J. M. Campbell, R. Frederix, F. Maltoni and F. Tramontano, Phys. Rev. Lett. **102**, 182003 (2009), [arXiv:0903.0005 [hep-ph]].
- [60] **SM and NLO Multileg Working Group** Collaboration, J. R. Andersen *et al.*, arXiv:1003.1241 [hep-ph].
- [61] Z. Sullivan, Phys. Rev. **D70**, 114012 (2004), [hep-ph/0408049].
- [62] **CDF** Collaboration, J. Lueck, Private Communication .
- [63] S. Frixione and B. R. Webber, JHEP **06**, 029 (2002), [hep-ph/0204244].
- [64] P. Nason, JHEP **11**, 040 (2004), [hep-ph/0409146].